

Connecting vacua
of half-maximal supergravities :
a type IIA example

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with G. Dibitetto and D. Roest : [arXiv:1102.0239](https://arxiv.org/abs/1102.0239)

The footprint of extra dimensions

- › Four dimensional supergravity theories appear when compactifying string theory
- › Fluctuations of the internal space around a fixed geometry translates into **massless** 4d **scalar fields** known as “*moduli*”

$$\mathcal{L} = \frac{1}{2} R - \frac{1}{2} \partial_\mu \phi_i \partial^\mu \phi^i$$



Deviations
from GR !!

massless scalars = long range interactions (precision tests of GR)

Linking strings to observations  Mechanisms to stabilise moduli !!

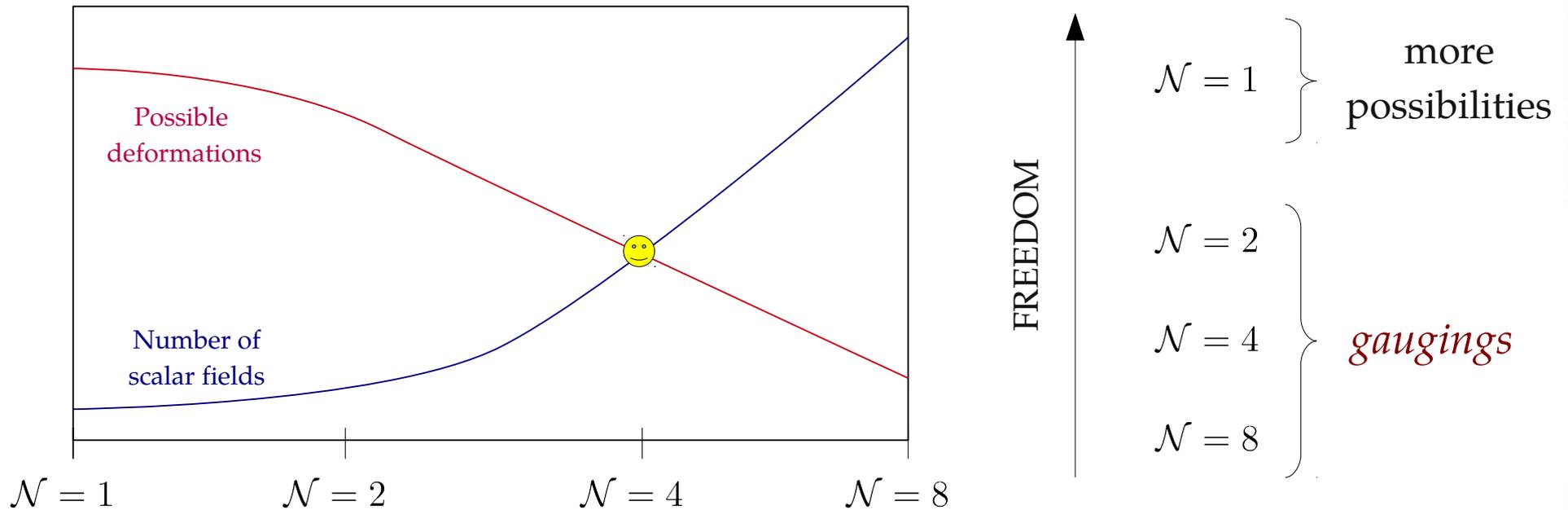
$$V(\phi) = m_{ij}^2 \phi^i \phi^j + \dots$$

- › Moduli VEVs $\langle \phi \rangle = \phi_0$ determine 4d physics

$\Lambda_{c.c} \equiv V(\phi_0)$
 g_s and Vol_{int}
fermi masses

How to deform massless theories to have $V(\phi) \neq 0$?

- Supersymmetry dictates what deformations are allowed



gaugings = part of the global symmetry is promoted to local (*gauge*)

Questions :

- Can the whole vacuum structure be charted in $\mathcal{N} = 4$ theories ?
- Are there connections in the landscape of vacua ?

Half-maximal supergravities

‣ $\mathcal{N} = 4$ supergravities are commonly found in string reductions

‣ Global symmetry group

$$G = SL(2) \times SO(6, 6)$$

‣ **Field content** = supergravity multiplet + six vector multiplets

‣ **Vectors** $A_{\mu}^{\alpha M}$ in the fundamental of G

$\alpha = +, -$ is an electric-magnetic $SL(2)$ index
 $M = 1, \dots, 12$ is an $SO(6, 6)$ index } **24 vectors**

‣ The **scalar sector** parameterises the coset space $\mathcal{M} = G/H$ where H is the maximal compact subgroup of G

- 1 axion + 1 dilaton in $SL(2)$
 - 30 axions + 6 dilatons in $SO(6, 6)$
- } **38 scalars**

Gaugings and scalar potential

[Schon, Weidner '06]

‣ A subgroup $G_0 \subset SL(2) \times SO(6, 6)$ is promoted to local (*gauged*)

‣ *Gaugings* are classified by the **embedding tensor parameters**

$$\xi_{\alpha M} \in (\mathbf{2}, \mathbf{12}) \quad \text{and} \quad f_{\alpha MNP} \in (\mathbf{2}, \mathbf{220})$$

‣ Supersymmetry + gauge invariance determine the **scalar potential**

$$V(\Phi) = \sum_{\text{terms}} f f \Phi^{\text{high degree}} + \xi \xi \Phi^{\text{high degree}}$$

Quadratic in the emb. tens. parameters !!

The SO(3) truncation

› Keeping SO(3)-invariant fields and embedding tensor parameters

- global symmetry $G = SL(2) \times SO(2, 2)$
- $\xi_{\alpha M} = 0$
- scalar coset = **3 complex scalars** = *STU* - models !!

[Derendinger, Kounnas, Petropoulos, Zwirner '04]

› The *gaugings* $G_0 \subset G$ and the scalar potential $V(S, T, U)$ are specified by the embedding tensor parameters $f_{\alpha MNP}$

gaugings

$$A_\mu = A_\mu^{\alpha M} T_{\alpha M}$$

$$[T_{\alpha M}, T_{\beta N}] = f_{\alpha MN}{}^P T_{\beta P}$$



Quadratic Constraints

$$\epsilon^{\alpha\beta} f_{\alpha MNR} f_{\beta PQ}{}^R = 0$$

$$f_{\alpha R[MN} f_{\beta PQ]}{}^R = 0$$

› String embedding as type II orientifold reductions

$$\text{generalised fluxes} = f_{\alpha MNP}$$

[Dibitetto, Linares, Roest '10]

We would like to . . .

- 1) build **all** the consistent SO(3)-invariant gaugings specified by $f_{\alpha MNP}$ by solving the quadratic constraints

$$\epsilon f f = 0 \quad \text{and} \quad f f = 0$$

- 2) compute **all** the SO(3)-invariant extrema of the f -induced scalar potential $V(f, \Phi)$ by solving the extremisation conditions

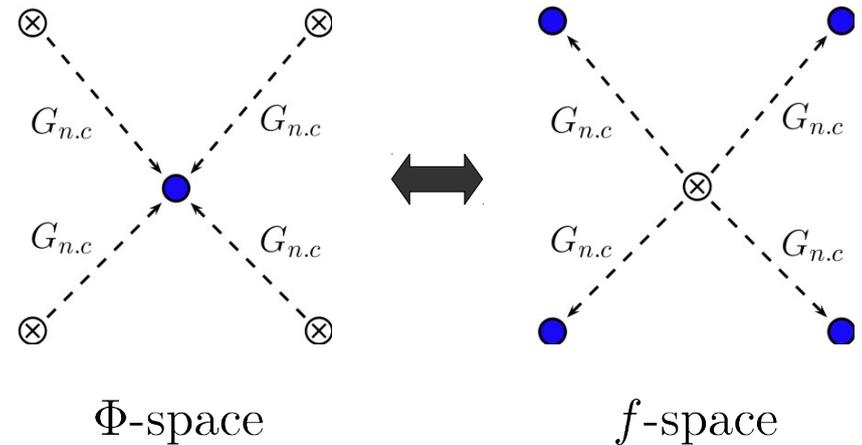
$$\left. \frac{\partial V}{\partial \Phi} \right|_{\Phi_0} = 0 \quad \text{with} \quad \Phi \equiv (S, T, U)$$

- 3) check stability of these extrema with respect to fluctuations of **all** the 38 scalars of half-maximal supergravity
- 4) identify the gauge group G_0 underlying **all** the different solutions

. . . but is this doable ?

Strategy and tools

- **Idea** : use the global symmetry group (non-compact part) to bring **any** field solution back to the origin !!



- At the origin everything is **simply quadratic** in the $f_{\alpha MNP}$ parameters

computing the vacua structure	=	solving a quadratic ideal $I = \langle \partial_{\Phi} V _{\Phi_0} , \epsilon f f , f f \rangle$
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- Algebraic Geometry algorithms : GTZ prime decomposition , ...

[Singular project, 97]

$I = J_1 \cap J_2 \cap \dots \cap J_n$	➔	Splitting of the landscape into n disconnected pieces !!
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An example : type IIA with metric fluxes

- › Testing the method with type IIA orientifold models including **gauge fluxes** and a **metric flux**

[Dall'Agata, Villadoro, Zwirner '09]

$$\left(F_{p=0,2,4,6} , H_3 \right) + \omega \subset f_{\alpha MNP}$$

Q.C. of *gaugings* = B.I. + tadpoles cancellation

- › **Subset** of embedding tensor components **closed under** $G_{n,c}$
 - ✓ Fields can still be set at the origin without lost of generality
 - ✓ Stability with respect to fluctuations around the origin can be computed

[Borghese, Roest '10]

- › **Vacua structure** of these type IIA orientifolds



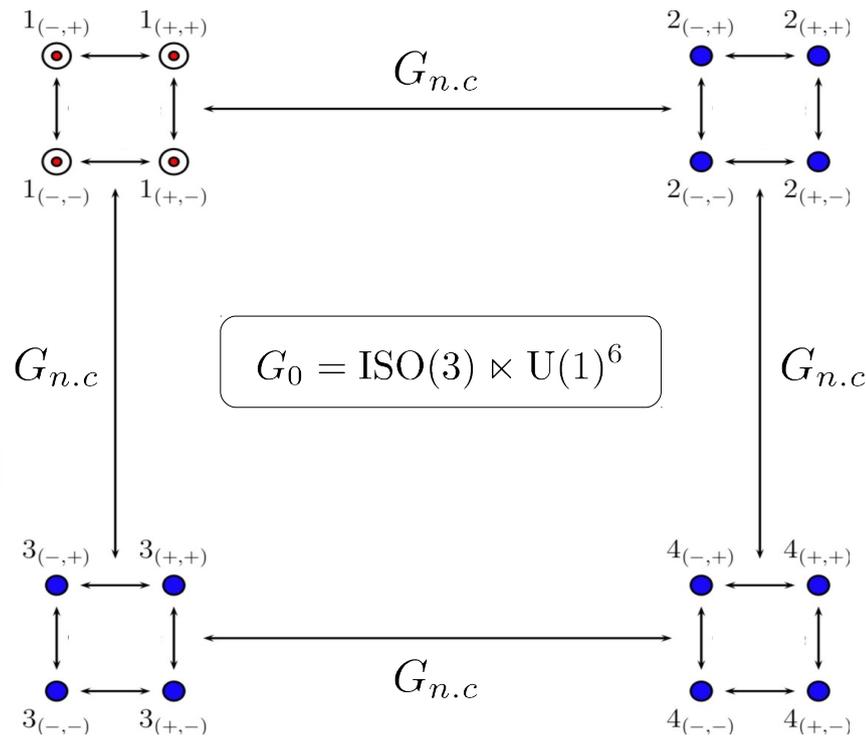
The 16 critical points

➤ An **AdS₄** landscape

$$16 = 4 + 4 + 4 + 4$$

➤ **All** the solutions are embeddable in $\mathcal{N} = 8$

$1_{(\pm,\pm)}$	$2_{(\pm,\pm)}$	$3_{(\pm,\pm)}$	$4_{(\pm,\pm)}$
$\mathcal{N} = 1$ SUSY & FAKE SUSY	SUSY	SUSY	SUSY
stable	unstable	stable	stable
$m^2 = -2/3$	$m^2 = -4/5$	$m^2 > 0$	$m^2 > 0$
$V = -1$	$V = -32/27$	$V = -8/15$	$V = -32/27$



(*) $m^2 \equiv$ lightest mode (B.F. bound = $-3/4$)

- **All** the solutions are connected
- **Unique gauging**

Unique theory
with **4 different vacua !!**

Conclusions

- Some progress towards disentangling the landscape of half-maximal supergravities can still be done without performing statistics of vacua
- The approach relies on the combined use of global symmetries and of algebraic geometry techniques
- As a warming-up, the complete vacua structure of simple type IIA orientifold theories can be worked out revealing some odd features :
 - i)* stability without supersymmetry
 - ii)* connections between vacua
 - iii)* $\mathcal{N} = 8$ embedding of the entire vacua structure
- **For the future :**
 - Going beyond the geometric limit : non-geometric backgrounds . . .
 - Systematic search of de Sitter stable solutions in extended supergravity and also links to Cosmology

Thanks for your attention !!

Extra material...

Internal geometries and massless theories . . .

“maximal”

$$\mathbb{T}^6 = \text{torus} \times \text{torus} \times \text{torus}$$

$$\mathcal{N} = 8 \quad (70 \text{ scalars})$$

“minimally extended”

$$CY_3 = \text{Calabi-Yau 3-fold}$$

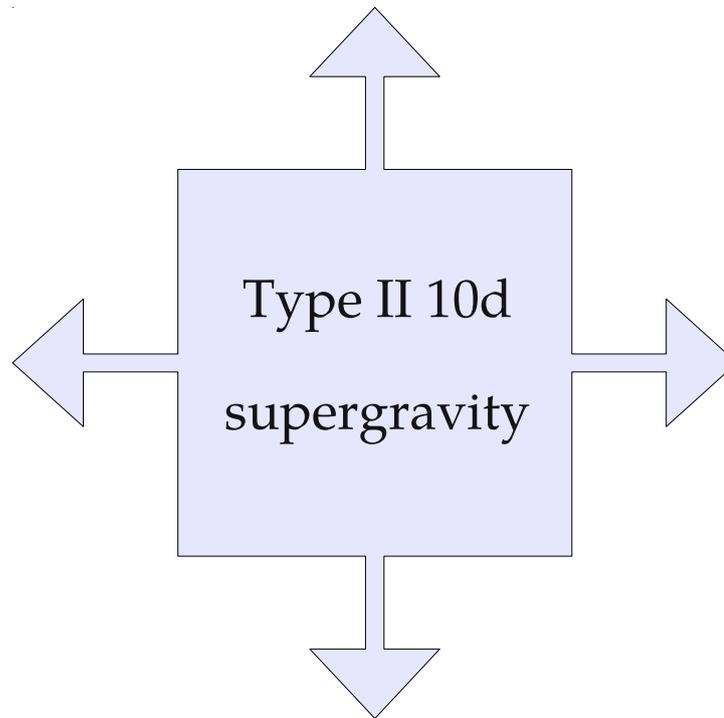
$$\mathcal{N} = 2$$

$$\text{scalars} \leftrightarrow (h^{(1,1)}, h^{(1,2)})$$

Orientifolds of \mathbb{T}^6

$$\mathcal{N} = 4 \quad (38 \text{ scalars})$$

“half-maximal”



Orientifolds of CY_3

$$\mathcal{N} = 1 \quad \text{scalars} \leftrightarrow (h^{(1,1)}, h^{(1,2)})$$

“minimal”

Lower bound on topologically distinct CY_3
30108

Gaugings and their higher-dimensional origin

- Scalars potentials are induced by “*gaugings*”: Part of the global symmetry is promoted to local (*gauge*)

[de Wit, Samtleben, Trigiante '07]

[Schon, Weidner '06]

- $\mathcal{N} = 8$: Gauging a subgroup of the global symmetry $G = E_7$

Internal space extension \longleftrightarrow Exceptional Generalised Geometry ?

[Pacheco, Waldram '08 , Grana, Louis, Sim, Waldram '09]

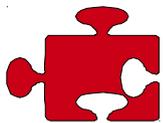
[Aldazabal, Andrés, Cámara, Grana '10]

- $\mathcal{N} = 4$: Gauging a subgroup of the global symmetry $G = SL(2) \times SO(6, 6)$

Internal space extension \longleftrightarrow Doubled/Generalised Geometry ?

[Hitchin '02, Gualtieri '04]

[Hull '04, '06]



pathological
internal spaces

String compactifications including
generalised flux backgrounds !!

De Sitter in extended supergravity

➤ $\mathcal{N} = 8$: **unstable** dS solutions with $SO(4,4)$ and $SO(5,3)$ gaugings
[Hull, Warner '85]

➤ $\mathcal{N} = 4$: **unstable** dS solutions with gaugings at *angles*
[De Roo, Wagemans '85]

i) $G_1 \times G_2$ gaugings with $\left\{ \begin{array}{l} G_i = SO(p_i, q_i) \quad , \quad p_i + q_i = 4 \\ G_i = CSO(p_i, q_i, r_i) \quad , \quad p_i + q_i + r_i = 4 \end{array} \right.$

[De Roo, Westra, Panda, (Trigiante) '02, '03, '06]

ii) $SO(3,1) \times U(1)^6$ gauging

[Dibitetto, A.G, Roest '11]

non-geometric fluxes in string theory !!

[Dibitetto, Linares, Roest '10]

➤ $\mathcal{N} = 2$: **stable** dS solutions with $SO(2,1) \times SO(3)$ gauging plus Fayet-Iliopoulos terms
[Fré, Trigiante, Van Proeyen '03]

unclear origin in string theory !!

De Sitter in minimal supergravity

- No-go theorems forbidding dS solutions in $\mathcal{N} = 1$ compactifications with magnetic fluxes

$$V_o = -\frac{1}{9} \sum \bar{F}^2 \leq 0 \quad \Rightarrow \quad \text{AdS !!}$$

[Hertzberg, Kachru, Taylor, Tegmark '07]

- Including **more general fluxes** : (metric + non-geometric)

$$V_o = -\frac{1}{9} \sum \bar{F}^2 + \Delta V_{\text{metric}} + \Delta V_{\text{non-geom}}$$

- a) metric fluxes \longleftrightarrow **unstable** dS in type IIA models

[Caviezel, Koerber, Kors, Lust, Wrase, Zagerman '08]

- b) non-geometric fluxes \longleftrightarrow **stable** dS in type IIA models

[de Carlos, A.G, Moreno '09, '10]

- Including D-branes to **uplift an AdS** solution

[Kachru, Kallosh, Linde, Trivedi '03]

- a) D-terms from D-branes \longleftrightarrow **stable** dS in type IIB models

[Burgess, Kallosh, Quevedo '03]

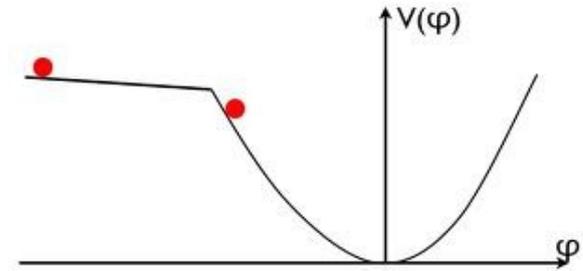
- b) non-perturbative effects from D-branes \longleftrightarrow **stable** dS in type IIB

[Achúcarro, de Carlos, Casas, Doplicher '06]

Cosmology from moduli ?

› slow-roll inflation requires an almost flat dS saddle point of $V(\phi)$ from which to start rolling down

$$\eta \equiv M_p^2 \left(\frac{V''}{V} \right) \ll 1$$



› dS saddle points **suffering from eta-problem**, *i.e.* $\eta \sim \mathcal{O}(1)$

i) gaugings in extended supergravity

[Kallosh, Linde, Prokushkin, Shmakova '01]

ii) general fluxes in minimal supergravity

[Flauger, Paban, Robbins, Wrase '08]

[de Carlos, A.G, Moreno '10]

› dS saddle points with $\eta \ll 1$ in minimal supergravity including non-perturbative effects \Rightarrow **axion inflation !!**

[Dimopoulos, Kachru, McGreevy, Wacker '05]